# GEnERALIZED LORENTZ Transformations 

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This article presents the generalized Lorentz transformations of time, space, velocity and acceleration which can be applied in any inertial or non-inertial ( uniform circular motion) frame.

## Introduction

If we consider an inertial or non-inertial ( uniform circular motion ) frame $S$ and another inertial frame $\Sigma$ whose origins coincide at time zero (in both frames ) then the time $(t)$, the position $(\mathbf{r})$, the velocity $(\mathbf{v})$ and the acceleration (a) of a ( massive or non-massive ) particle relative to the inertial frame $\Sigma$ are:

$$
\begin{aligned}
& t=\int_{0}^{\mathrm{t}} \gamma \mathrm{dt}-\gamma \frac{\vec{r} \cdot \mathbf{V}}{c^{2}} \\
& \mathbf{r}=\vec{r}+\frac{\gamma^{2}}{\gamma+1} \frac{(\vec{r} \cdot \mathbf{V}) \mathbf{V}}{c^{2}}-\gamma \mathbf{R}-\frac{\gamma^{2}}{\gamma+1} \frac{(\mathbf{R} \times \mathbf{V}) \times \mathbf{V}}{c^{2}} \\
& \mathbf{v}=\frac{d \mathbf{r}}{d t} \\
& \mathbf{a}=\frac{d \mathbf{v}}{d t}
\end{aligned}
$$

where $(\mathrm{t}, \vec{r})$ are the time and the position of the particle relative to the frame S , $(\mathbf{R}, \mathbf{V}, \mathbf{A})$ are the position, the velocity and the acceleration of the origin of the inertial frame $\Sigma$ relative to the frame $\mathrm{S},(c)$ is the speed of light in vacuum, and $\gamma=\left(1-\mathbf{V} \cdot \mathbf{V} / c^{2}\right)^{-1 / 2}$

- $\frac{d \mathbf{r}}{d t}=\left(\frac{d \mathbf{r}}{d \mathrm{t}}+\Omega \times \mathbf{r}\right)\left(\frac{1}{d t / \mathrm{dt}}\right)$
- $\frac{d \mathbf{v}}{d t}=\left(\frac{d \mathbf{v}}{\mathrm{dt}}+\Omega \times \mathbf{v}\right)\left(\frac{1}{d t / \mathrm{dt}}\right)$
- $\Omega=\frac{\gamma^{2}}{\gamma+1} \frac{(\mathbf{A} \times \mathbf{V})}{c^{2}}$
- $\frac{\gamma^{2}}{\gamma+1} \frac{1}{c^{2}}=\frac{\gamma-1}{\mathbf{V}^{2}} \quad\left(\mathbf{V}^{2}=\mathbf{V} \cdot \mathbf{V}\right)$
- $\gamma \mathbf{R}+\frac{\gamma^{2}}{\gamma+1} \frac{(\mathbf{R} \times \mathbf{V}) \times \mathbf{V}}{c^{2}}=\mathbf{R}+\frac{\gamma^{2}}{\gamma+1} \frac{(\mathbf{R} \cdot \mathbf{V}) \mathbf{V}}{c^{2}}$


## General Observations

If the frame $S$ is inertial then $(\mathbf{A}=0),(\mathbf{V}=$ cte $),(\mathbf{R}=\mathbf{V} \mathrm{t}),(\gamma=$ cte $)$ $\left(\int_{0}^{\mathrm{t}} \gamma \mathrm{dt}=\gamma \mathrm{t}\right),(\mathbf{R} \times \mathbf{V}=0) \&(\Omega=0)$

If the frame $\mathbf{S}$ is non-inertial (uniform circular motion) then ( $\mathbf{A} \neq 0$ ) $(\mathbf{A} \cdot \mathbf{V}=0),(\mathbf{V} \cdot \mathbf{V}=$ cte $),(\gamma=$ cte $)\left(\int_{0}^{\mathrm{t}} \gamma \mathrm{dt}=\gamma \mathrm{t}\right) \&(\Omega \neq 0)$

In addition, if the frame $S$ is non-inertial ( uniform circular motion) then the observer $S$ should preferably use an origin $\mathrm{O}^{\prime}$ such that $(\mathbf{R} \cdot \mathbf{V}=0)$

## Bibliography

[1] R. A. Nelson, J. Math. Phys. 28, 2379 (1987).
[2] R. A. Nelson, J. Math. Phys. 35, 6224 (1994).
[3] C. Møller, The Theory of Relativity (1952).

